

Supplementary information

**Few-fs resolution of a photoactive protein
traversing a conical intersection**

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1 **Supplementary Information**

2 **Few-fs resolution of a photoactive protein traversing a conical intersection**

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35 1. Data-analytical approach

36 Our approach is based on manifold-based machine learning, including Nonlinear Laplacian
37 Spectral Analysis¹⁹. In this approach, data vectors are ordered based on their known
38 timestamps, and concatenated to form the supervector matrix X . The supervectors are then
39 projected onto their manifold,

$$40 \text{ viz. } A = X\mu\Phi. \quad [1]$$

41 Here, μ and Φ are respectively the Riemannian measure and the Diffusion Map empirical
42 orthogonal functions (EOF).

$$43 \text{ Singular Value Decomposition: } A = USV^T \quad [2]$$

$$44 \text{ and back projection: } \tilde{X} = A\Phi^T = USV^T\Phi^T \quad [3]$$

45 are applied to yield the reconstruction matrix \tilde{X} , which must be unwrapped to give individual
46 reconstructed data vectors^{16,19}.

47

48 Independent orthogonal dynamical modes can be studied by reconstructing with specific SVD

$$49 \text{ modes: } \tilde{X}_1 = U_1S_1V_1^T\Phi^T, \tilde{X}_2 = U_2S_2V_2^T\Phi^T, \dots, \tilde{X}_k = U_kS_kV_k^T\Phi^T. \quad [4]$$

50

51 2. Computing Euclidean distances and dot products

52 For N data vectors with D pixels each, and concatenation parameter c , the supervector matrix X
53 (dimensions $cD \times (N - c + 1)$) can be huge, even for modest values of N and D . It is, however,
54 not necessary to explicitly store or manipulate X . For instance, the SVD step above (Equation
55 [2]) can be more efficiently carried out by using the following steps:

56 (i) calculate the dot products amongst the supervectors, i.e. $X^T X$, in blocks (more details below)

$$57 \text{ (ii) form the } A^T A \text{ matrix, i.e. } A^T A = (\mu\Phi)^T X^T X (\mu\Phi); \quad [5]$$

- 58 (iii) solve for the eigenvalues and eigenvectors of the $A^T A$ matrix;
- 59 (iv) the right singular vectors (V) of A are the eigenvectors of $A^T A$; and the singular values (S) of
- 60 A are the square roots of the eigenvalues of $A^T A$, or in other words,
- 61 $A^T A = (\mu\Phi)^T (X^T X) \mu\Phi = VS^2V^T$; [6]
- 62 (v) the left singular vectors (U) of A are obtained from $U = X\mu\Phi VS^{-1}$. [7]

63

64 Note that if we are using a small number ℓ_{max} of Diffusion Map EOFs, say $\ell_{max} = 100$, the

65 matrix $A^T A$, of dimensions $\ell_{max} \times \ell_{max} = 100 \times 100$, is rather small, and can be accumulated

66 using a double loop through the block structure of $X^T X$. Also, since a full reconstruction results

67 in up to c copies of each individual snapshot, which might be too many, it is not necessary to

68 calculate the full U matrix. Equation [7] can thus be used to compute U in a row-wise/ block-

69 wise fashion to only generate enough copies of each individual snapshot for our reconstruction.

70

71 Squared Euclidean distances and dot products amongst supervectors are calculated in Nonlinear

72 Laplacian Spectral Analysis (NLSA). For N data vectors with D pixels each, and concatenation

73 parameter c , runtimes for these steps scale as $\sim N^2 c D$. Calculations with N , D , and c in the tens

74 or hundreds of thousands can, literally, take years on a desktop machine.

75

76 For this paper, we have developed a so-called Shift-and-Add algorithm, which reduces the

77 runtime scaling to $\sim N^2 D + N^2 \log_2(c)$. Calculations with N , D , and c in the tens or hundreds of

78 thousands now take only days on a desktop machine, and only hours on computer clusters with

79 fairly modest resources. To describe this algorithm in more detail, we define:

80 $\vec{x}_j =$ data vector j ,

81 \vec{x}_j^c = supervector j with concatenation parameter c , and

82 $s_{i,j}^c$ = the squared Euclidean distance between supervectors i and j . [8]

83 By definition, we have:

84 $s_{i,j}^c = |\vec{x}_i^c - \vec{x}_j^c|^2$. [9]

85

86 Writing out the constituent data vectors of the supervectors explicitly, Equation [9] becomes:

87 $s_{i,j}^c = \sum_{p=0}^{p=c-1} |\vec{x}_{i+p} - \vec{x}_{j+p}|^2$. [10]

88

89 For concatenation parameter a , where $a < c$, we break up the sum in Equation [10] to give:

90 $s_{i,j}^c = \sum_{p=0}^{p=a-1} |\vec{x}_{i+p} - \vec{x}_{j+p}|^2 + \sum_{p=a}^{p=c-1} |\vec{x}_{i+p} - \vec{x}_{j+p}|^2$. [11]

91

92 Substituting $p = q + a$ in the second sum above yields:

93 $s_{i,j}^c = \sum_{p=0}^{p=a-1} |\vec{x}_{i+p} - \vec{x}_{j+p}|^2 + \sum_{q=0}^{q=c-a-1} |\vec{x}_{i+a+q} - \vec{x}_{j+a+q}|^2 = s_{i,j}^a + s_{i+a,j+a}^{c-a}$. [12]

94

95 Using Equation [12], the matrix of squared Euclidean distances amongst supervectors for any

96 concatenation parameter can be built from the matrices with lower concatenation parameters.

97 For example, starting with the matrix of squared Euclidean distances amongst data vectors, the

98 matrices of squared Euclidean distances between supervectors with concatenation parameters

99 $c=2$ and $c=4$ can be successively assembled as:

100 $s_{i,j}^{c=1} = |\vec{x}_i - \vec{x}_j|^2$,

101 $s_{i,j}^{c=2} = s_{i,j}^{c=1} + s_{i+1,j+1}^{c=1}$,

102 $s_{i,j}^{c=4} = s_{i,j}^{c=2} + s_{i+2,j+2}^{c=2}$. [13]

103 After the calculation of $s_{i,j}^{c=1}$, it takes $\log_2(c)$ steps of “doubling” ($\sim N^2$ additions each) to reach
104 the concatenation parameter c . Runtime thus scales as $\sim N^2 D + N^2 \log_2(c)$. [14]

105

106 Matrices of squared Euclidean distances amongst supervectors for arbitrary concatenation
107 parameters can be assembled, for instance

$$108 \quad s_{i,j}^{c=3} = s_{i,j}^{c=1} + s_{i+1,j+1}^{c=2} ,$$

$$109 \quad s_{i,j}^{c=5} = s_{i,j}^{c=2} + s_{i+2,j+2}^{c=3} ,$$

$$110 \quad s_{i,j}^{c=6} = s_{i,j}^{c=3} + s_{i+3,j+3}^{c=3} . \quad [15]$$

111

112 Starting with the elements of the matrix of squared Euclidean distances between data vectors in
113 files (blocks), the results for successively higher concatenation parameters can be obtained as
114 follows:

- 115 (i) Read files two at a time;
- 116 (ii) Shift the content of one with respect to the other; and
- 117 (iii) Add and save the results in files.

118 The above algorithm is named “Shift-and-Add”.

119

120 By replacing $|x - y|^2$ with $(x \cdot y)$ in the discussion above, it is obvious that Shift-and-Add can
121 be used to calculate the matrix of dot products amongst supervectors with arbitrary concatenation
122 parameters.

123

124

125

126 3. Modifications needed to handle sparse data matrices

127 Since our data matrix initially contains undefined elements (see Methods section entitled “Data
128 preprocessing”), we must adjust the way we calculate squared distances and projection in NLSA.

129 This adjustment is based on the number of times each unique reflection has been measured
130 (across the dataset), and by pre-normalizing (dividing) each row of the data matrix by the
131 number of times the corresponding Bragg reflection has been measured.

132

133 For squared distances:

- 134 i. Squared distance between two data vectors is calculated using only pixels defined in both
135 vectors;
- 136 ii. Squared distance between two data vectors with no common pixels is set to infinity, and
137 any supervector squared distance they contribute to will also be infinity (see section
138 above for the “Shift-and-Add” algorithm);
- 139 iii. Infinities in the squared distance matrix are removed/ignored in Diffusion Map where
140 only a small number of nearest-neighbor squared distances are kept.

141

142 To project on to the manifold ($X\mu\Phi = USV^T$) in NLSA:

- 143 i. Undefined pixels in the data matrix are set to 0;
- 144 ii. The dot-product $X^T X$ is calculated using Shift-and-Add (see section above for the “Shift-
145 and-Add” algorithm);
- 146 iii. V and S are obtained by solving for the eigenvectors/ eigenvalues of the matrix
147 $(\mu\Phi)^T(X^T X)\mu\Phi$, i.e. $(\mu\Phi)^T(X^T X)\mu\Phi = VS^2V^T$;
- 148 iv. U is obtained from $U = X\mu\Phi VS^{-1}$.

149 **4. Time-labeling of reconstructed videos**

150 In the time-lagged embedding used in this paper, the data vectors are ordered based on their
151 known timestamps, and concatenated to form the supervector matrix X . This matrix is then
152 projected onto its manifold Φ , and singular value decomposition and back projection are applied
153 to obtain the reconstructed matrix \tilde{X} in the data space, which must be unwrapped to give the
154 individual (reconstructed) data vectors^{16,19}.

155
156 When applied to data with inaccurately known timestamps, our data-analytical pipeline has been
157 shown to recover the dynamics on a uniform grid of timepoints with negligible timing error¹⁶.

158
159 Defining the timestamp of a supervector as the average of the timestamps of its constituent data
160 vectors, the concatenation parameter c is chosen so that:

- 161 i. The set of time steps Δt between consecutive supervectors in X becomes more or less
162 uniform; and
163 ii. The time step Δt between consecutive supervectors remain more or less constant as the
164 concatenation parameter is further increased.

165
166 For the present study, $c = 32768$, and $\Delta t = 7.35as$.

167
168 The columns of the reconstruction matrix \tilde{X} have the same supervector timestamps as the matrix
169 X , the individual constituent data vectors in \tilde{X} are, however, uniformly spaced with time step Δt .

170
171 The start time (t_{start}) of a reconstructed movie is determined by

- 172 i. Knowing the timestamp of the first supervector: \tilde{t}_1 ;
- 173 ii. Noting that the first data vector is half a concatenation window behind the supervector to
- 174 which it belongs: $-\frac{c-1}{2}\Delta t$;
- 175 iii. Knowing the number (p) of early data vectors that have been dropped, because they have
- 176 too few copies in the reconstruction: $+p\Delta t$.

177 Finally, $t_{start} = \tilde{t}_1 + \left(p - \frac{c-1}{2}\right)\Delta t$. [16]

178

179 For the results presented in this paper, $p = c = 32768$, and $\tilde{t}_1 = 164.24fs$. The start time of our

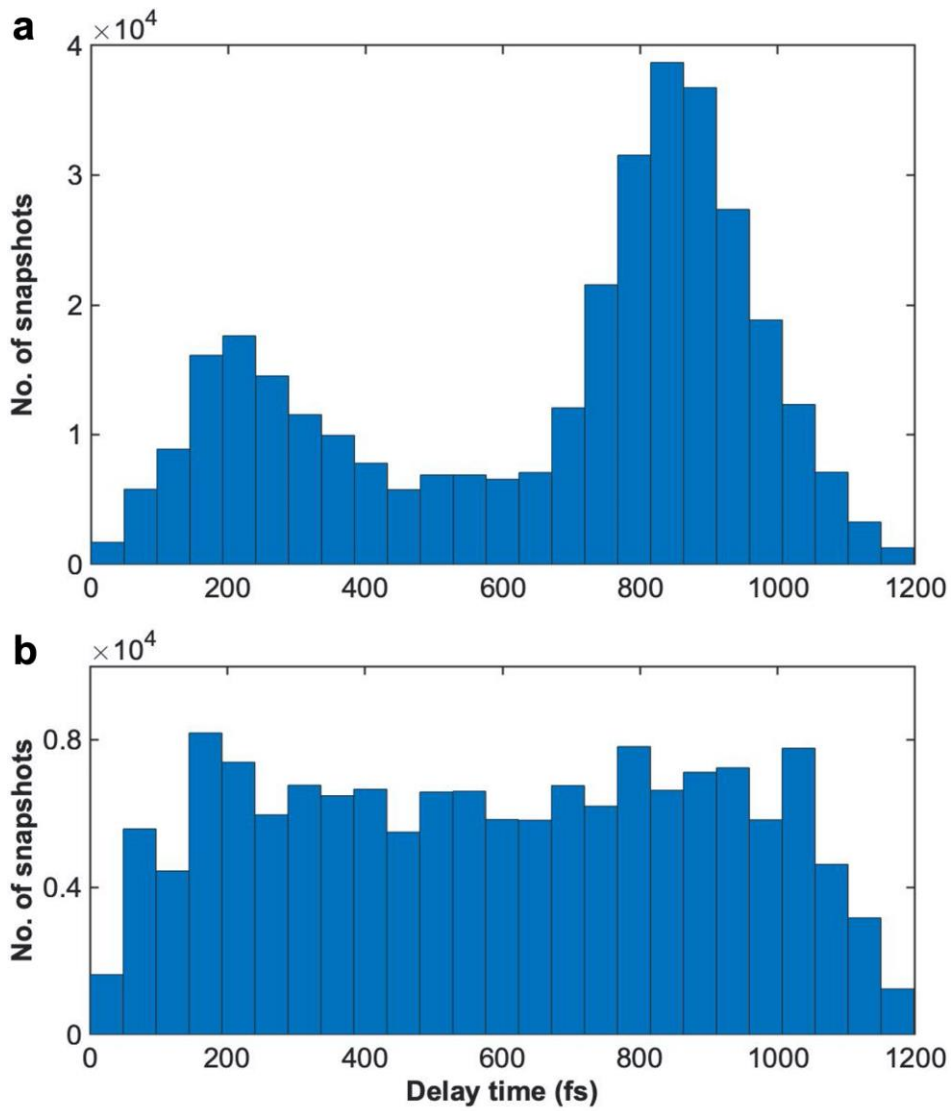
180 reconstructed movies is therefore:

181 $t_{start} = 164.24fs + \left(32768 - \frac{32768-1}{2}\right)7.35as = 284.67fs$. [17]

182

183

184



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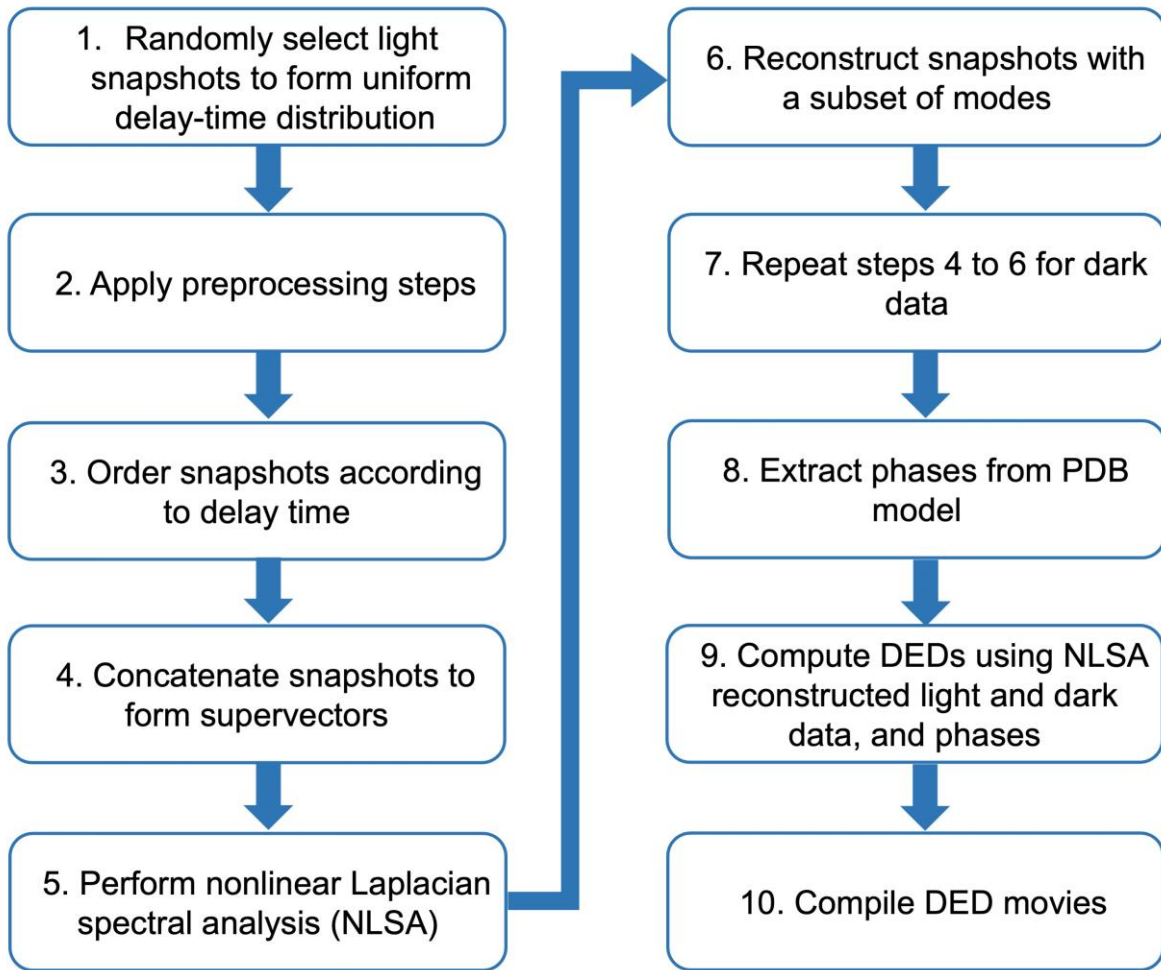
186 **Supplementary Fig. 1 | Histograms of the snapshot delay times. a,** Outcome of experiment.

187 **b,** After random subsampling of the experimental data to obtain a statistically uniform

188 distribution in delay time.

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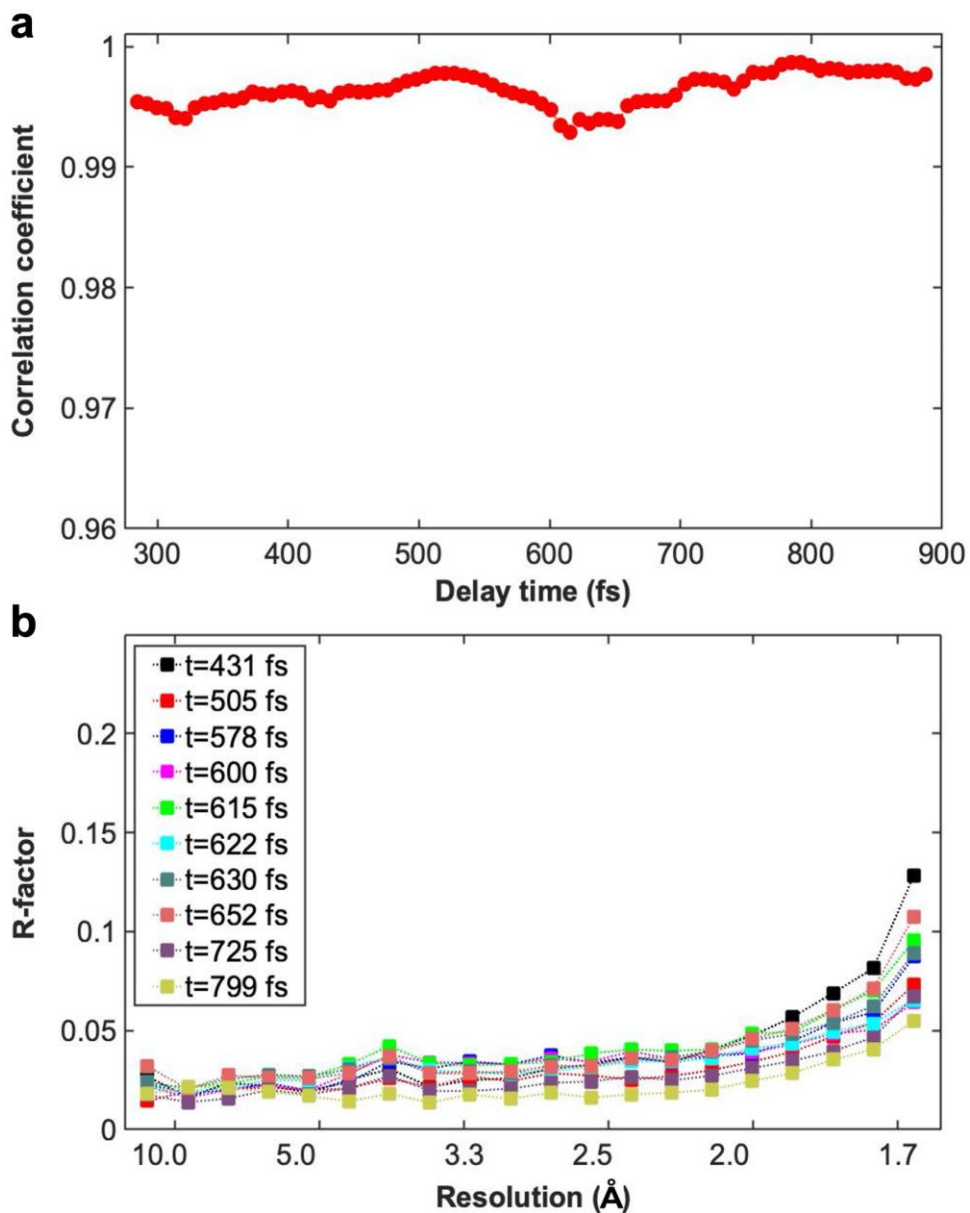
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Supplementary Fig. 2 | Flowchart of the analytical pipeline.



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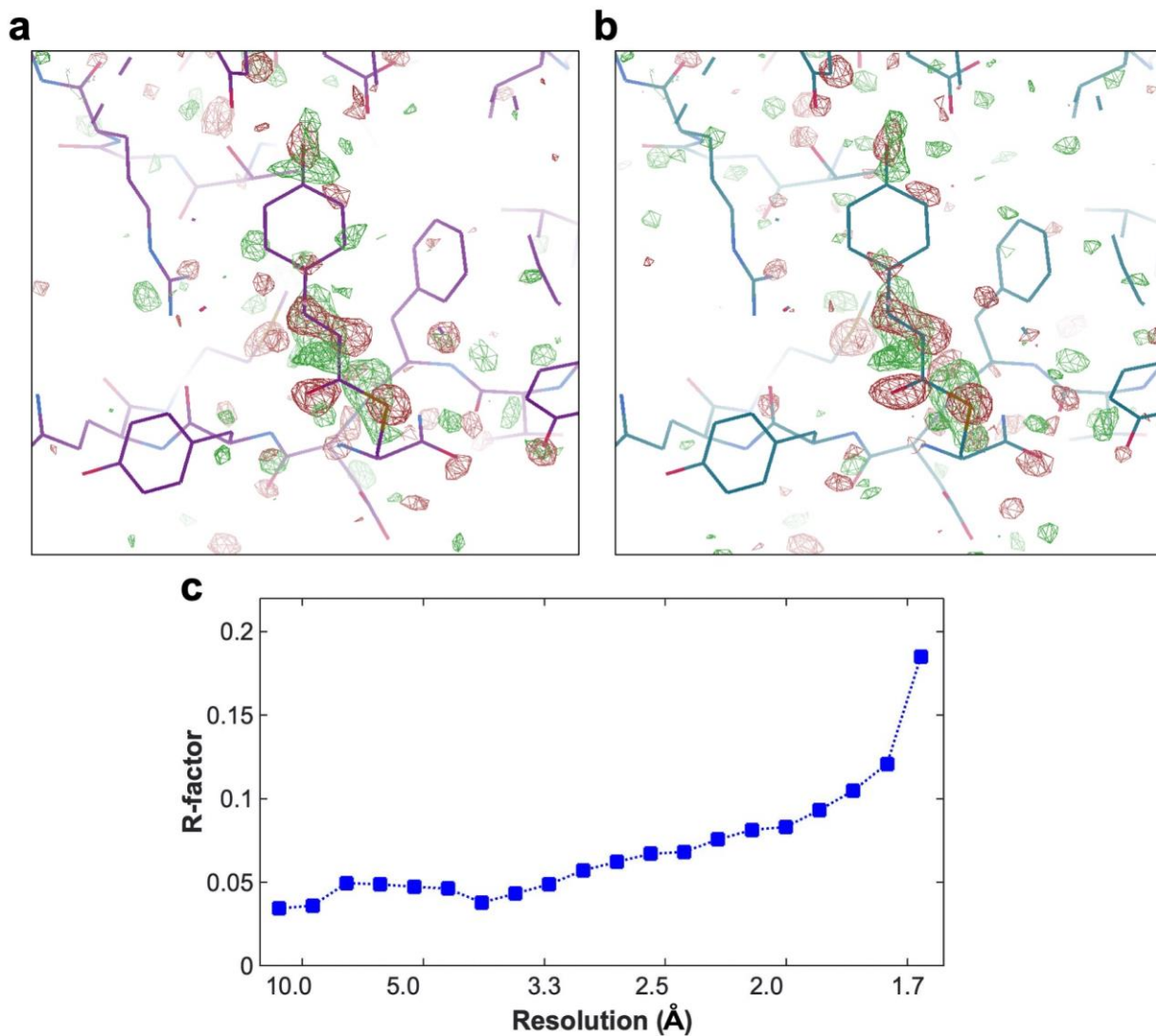
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199 **Supplementary Fig. 3 | Pearson correlation and R-factor between synthetic (input) and**
 200 **output diffraction volumes obtained from step 7 of Supplementary Fig. 2. a, Correlation.**

201 The average of correlation coefficients is 0.996. **b**, R- factor. Diffraction volumes in both cases
 202 were reconstructed using all non-noise NLSA modes.

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Supplementary Fig. 4 | Comparing difference electron density maps at 3 ps delay

207

obtained by: **a**, Standard time-resolved crystallographic analysis; **b**, Machine learning

208

algorithm used in this paper. Contour level for both maps: 3σ . **c**, R-factor between the

209

diffraction volumes at 3ps obtained by standard crystallographic approaches and that

210

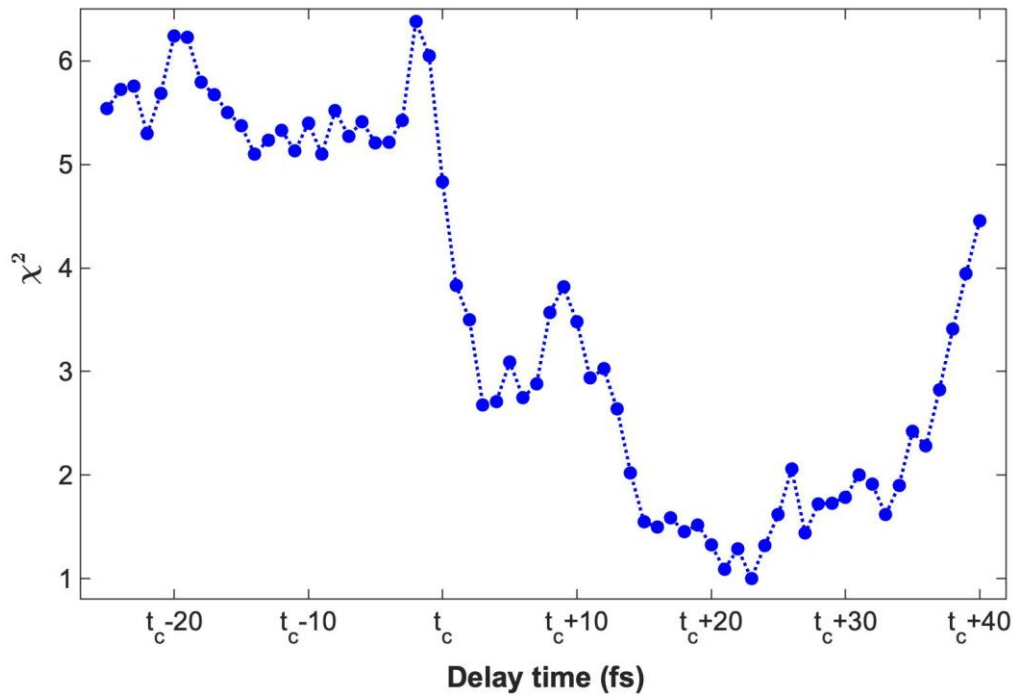
obtained by the analytical pipeline in this paper.

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217 **Supplementary Fig. 5** | χ^2 landscape of a typical best-fit, in this case for the mode2-mode5

218 combination, for different trajectory segments. The index t_c refers to the center of the 100-fs

219 timespan, which corresponds to the turning point in chronos.